

Exploring Non-Thermal Dark Matter through Monojet Events

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Abstract

Although the Standard Model (SM) successfully describes a wide range of particle interactions, it does not account for key phenomena such as dark matter and the matter—antimatter asymmetry. Various Beyond the Standard Model (BSM) scenarios have been proposed to address these limitations. This study investigates a BSM model featuring baryon-number-violating processes and a dark matter candidate. The model introduces two iso-singlet color-triplet scalars X at the TeV scale and one singlet Majorana fermion with a mass of approximately that of the proton. The fermion serves as the dark matter candidate, manifesting as missing transverse energy. The monojet channel is studied here.

Introduction

Standard Model of Elementary Particles

Event Selection

The monojet signature have light-jet, which is accompanied with MET and b-tagged jet

| Selection | | | |
|--------------------------------|-----------------------------|---|------------|
| | $N(u) \geq 1$ | • | U i b-t |
| Light quark jets Selection | $p_T(u) > 150 \mathrm{GeV}$ | | qu |
| | $ \eta(u) < 2.4$ | • | Le |
| | $N(b) \geq 1$ | | mu ron |
| B-tagged jets Selection | $p_T(b) > 20 \mathrm{GeV}$ | • | Ta |
| | $ \eta(b) < 2.4$ | | tau |

- U is used to denote a light jet
- b-tagged jets must be b-tagged jets and light quark jets must not be b-tagged
- Lepton Veto: Events containing electrons or muons with pT > 10 GeV and |eta| < 2.4 are removed.
- Tau-tagged Jet Veto: Events containing tau-tagged jets with pT > 20 GeV and |eta| < 2.4

- The Standard Model and BSM model
 - The Standard Model (SM) is a successful theory describing fundamental particles and forces, but the SM model has several limitations.
 - ▷ The need for Beyond the Standard Model (BSM) is presented
- Baryon asymmetry and existence of dark matter
- Imbalance between matter and antimatter in the universe,
 known as **baryon asymmetry**, which the SM model cannot
 fully explain
- Astronomical evidence supports the presence of dark matter, yet it remains undetected within the framework of the Standard Model

Theoretical Model Setup

- Baryon Number Violating Model (BNV)
- this model introduce TeV scale two-color-triplet scalars X_a and one singlet
 Majorana fermion, which is stable and can be candidate of dark-matter
- **These scalars** having hypercharge and existing color triplet SU(3)_c and color singlet SU(2)_L +4/3 include **interaction violating the baryon number** with $\Delta B=2$
- The minimal interaction Lagrangian can be expressed as :

 $\mathcal{L} \supset \lambda_{\alpha i} X_{\alpha} \psi u_{i}^{c} + \lambda_{\alpha i j}^{\prime} X_{\alpha}^{*} d_{i}^{c} d_{j}^{c} + \frac{m_{\psi}}{2} \bar{\psi}^{c} \psi + \text{h.c.} \ \bar{d}$

| the | a concrations of | matter | Interactions | I force corriere | |
|------------------------------|----------------------------|-----------------------------|-----------------------------|---------------------------|--|
| thr | (fermions) | (fermions) | | (bosons) | |
| I | II | III | | | |
| ass =2.16 MeV/c ² | =1.273 GeV/c ² | =172.57 GeV/c ² | 0 | =125.2 GeV/c ² | |
| rge 33 | ³⁵ C | ³⁵ t | 0 1 g | B H | |
| up | charm | top | gluon | higgs | |
| <u> </u> | | | | | |
| =4.7 MeV/c ² | =93.5 MeV/c ² | =4.183 GeV/c ² | 0 | | |
| 2 🦻 d | 15 S | 16 b | 1 Y | 9 | |
| down | strange | bottom | photon | 5 | |
| | | | Ċ | ' | |
| ≈0.511 MeV/c ² | ≈105.66 MeV/c ² | ≈1.77693 GeV/c ² | ≈91.188 GeV/c ² | | |
| ⁻¹ e | -1 ½ | -1 T | 0 Z | S S | |
| | | | | lō S | |
| electron | muon | tau | Z boson | Ss | |
| <0.8 eV/c ² | <0.17 MeV/c ² | <18.2 MeV/c ² | ≈80.3692 GeV/c ² | <u>a</u> 8 | |
| | ° V. | v Vr | ±1 1 | H ^m | |
| 2 76 | | 22 | ' 🖤 | ⊇ē | |
| electron | muon | T 311 | | | |



- The background was plotted separately for each sample.
- For the signals, **different mass points** with the same couplings were plotted.



- Monojet Signature
 - We are interested in monjet signature with
 production with X_a from fusion of d-b and s-b quarks
 only
 - For **monojet signature** studies, we focus on the fusion of *d-b* and *s-b* quarks and rewrite the relevant phenomenological Lagrangian as :
 - $\mathcal{L}_{\text{monojet}} \supset \lambda_{\psi u} X_1 \psi u^c + \lambda_{db} X_1^* d^c b^c + \lambda_{sb} X_1^* s^c b^c + \text{h.c.} \text{ the other fuse to produce the TeV-scale scalar } X_1, \text{ which}$
 - $\circ~$ Couplings λ_{db} and λ_{sb} are denoted $~\lambda_1$, and $\lambda_{\psi u}$ is renamed λ_2

Figure: Feynman diagram for the monojet signature: a d-quark or s-quark from one proton and a bottom-quark from the other fuse to produce the TeV-scale scalar X_1 , which subsequently decays into the Majorana dark matter fermion ψ and an up antiquark. The up antiquark hadronizes into the observed light jet, while the invisible ψ escapes the detector, yielding MET.

X

MC Sample Production

• Simulation setting

- MadGraph5 aMC@NLO (v2.6.7) : Simulation for proton proton collisions with $\sqrt{s} = 13$ TeV at LHC
- **PYTHIA8 (v8.244) :** The parton showering and hadronization simulation
- **DELPHES (v3.4.2) :** CMS Detector simulation
- Signal Sample Production
 - For TeV scale scalar X, mass of X1 (MX1) is set to 1, 1.5 and 2 TeV

| | MX1 = 1.0 T | ſeV | | | | | | | |
|-------------|---|--|---|--|--|--|--|--|--|
| - | | Cro | oss Section [| pb] | | | | | |
| - | 1.0 | 3.2952E-02 | 5.0999E-02 | 6.8389E-01 | 1.1593E+00 | | | | |
| | 0.5 | 3.1077E-02 | 4.6459E-02 | 2.6825E-01 | 3.4055E-01 | | | | |
| | 0.1 | 1.1135E-02 | 1.1688E-02 | 1.3038E-02 | 1.4239E-02 | | | | |
| - | 0.08 | 8.0256E-03 | 8.2541E-03 | 8.2671E-03 | 9.1387E-03 | | | | |
| _ | λ2 \ λ1 | 0.08 | 0.1 | 0.5 | 1.0 | | | | |
| λ_2 | • | MX1 = 1.0 TeV | | | | | | | |
| 1.0 | 25000 20000 15000 5000 5000 0 2000 400 600 800 1001 m | 25000 - 15000 - 15000 - 15000 - 5000 - - - - - - - - - - - - - - | 1400- 12000- 10000- 4000- 2000- 2000- 2000- 2000- 2000- 00-200-400 fe0- 2000- 00-200-400 fe0- 2000- | 5000 500 5000 5 | 0 400 500 100 1001 1001 1001 1001 1000 1000 mass of MX1 (GeV) | | | | |
| 0.5 | 60000- 50000- 30000- 30000- 20000- 10000- 0 200 400 660 860 10001 mm | 60000 - 50000 - 40000 - 30000 - 10000 - 10000 - 10000 - 10000 - 10000 - 10000 - 10000 - 10000 - 10000 - | 2000 2000 18000 19000 10000 10000 10000 2000 0001201400160016002000 0 200 400 60016002000 | 00 800 id001200 id00 id00 id00 id00 id00 id00 id00 i | And the second s | | | | |
| 0.1 | 90000 | 90000 - 80000 - 70000 - 90000 - 90000 - 90000 - 90000 - 10000 - 10000 - 10000 - 10000 - 10000 - 10000 - 10000 - 1000 - 10 | 25000 20000 15000 5000 5000 5000 5000 50 | 00 100013001400140014002500 mass of MX1 [GeV] | 0 400 600 100 100 100 100 100 100 100 100 1 | | | | |
| 0.08 | 80000 | 90006 80006 - 70006 - 60006 - 60006 - 4006 - 30006 - 30006 - 30006 - 30006 - 30006 - 30006 - 30006 - - 30006 - - - - - - - - - - - - - | 2000 2000 1500 10000 5000 5000 6 200 400 6 0 200 400 6 | 00 600 100012001400110012000 mass of MX1 [GeV] | 440 600 MOCI LOCI 1001 MOCI MOCI MOCI MOCI MOCI MOCI MOCI MOCI | | | | |
| | 0.08 0.1 0.5 1. | | | | | | | | |
| | Process Cross-section [pb] | | | | | | | | |
| | ר + up | Tbar to 3jets | 5.5 | 5046E+0 | 2 | | | | |
| | + up | W to 3jets | 4.6 | 6546E+0 | 4 | | | | |
| | Dr pr + up | ell-Yan ocess to 3jets | 8.5 | 5844E+0 | 3 | | | | |
| | V E | lector Boson | 4.3 | 3912E+0 | 1 | | | | |

- The figure on the left shows the transverse mass distribution of the MET and light quark jets, and right plot shows $\Delta \phi$ of MET and light quark.
- Above plots is few of the distributions used to train the BDT, with differences that help distinguish the signal from the background at each point.

BDT Evaluation



- The mass of X2 (MX2) is fixed by 10 TeV
- Signal will be generated for grid of λ_1 and λ_2 : {0.08, 0.1, 0.5, 1.0}
- Background Production
 - TTbar + on to 3 jets: Simulated in bins of scalar HT (hadronic transverse energy)
 - W+jets / Z+jets (Drell–Yan): Also generated in
 HT-binned format to cover wide kinematic regions efficiently.
 - Diboson (WW, WZ and ZZ): Inclusive generation was applied for fully leptonic and semileptonic decays, and HT-binned samples were used for specific hadronic modes (e.g., WW→lvqq).

- channel for the cut as a function of BDT value are shown in the plot on the right.
- In this case, the BDT value of 0.131 shows a maximum of the significance of 5.332.
 Results for other MX1s and coupling constants were also calculated.

Summary and Plan

- We study a BNV model featuring a DM candidate, using monojet events. A BDT is used here for optimization of sensitivity.
- To refine the exclusion limit using this monojet channel, we plan to generate additional samples across a danser coupling constant grid. Finally, more background samples will be generated for statistical improvement.

Reference

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